## Diffusive capture process on complex networks

Sungmin Lee, Soon-Hyung Yook,\* and Yup Kim<sup>†</sup>

Department of Physics and Research Institute for Basic Sciences, Kyung Hee University, Seoul 130-701, Korea (Descined 28 March 2006, sublished 20 October 2006)

(Received 28 March 2006; published 30 October 2006)

We study the dynamical properties of a diffusing lamb captured by a diffusing lion on the complex networks with various sizes of N. We find that the lifetime  $\langle T \rangle$  of a lamb scales as  $\langle T \rangle \sim N$  and the survival probability  $S(N \rightarrow \infty, t)$  becomes finite on scale-free networks with degree exponent  $\gamma > 3$ . However, S(N, t) for  $\gamma < 3$  has a long-living tail on tree-structured scale-free networks and decays exponentially on looped scale-free networks. This suggests that the second moment of degree distribution  $\langle k^2 \rangle$  is the relevant factor for the dynamical properties in the diffusive capture process. We numerically find that the normalized number of capture events at a node with degree k, n(k), decreases as  $n(k) \sim k^{-\sigma}$ . When  $\gamma < 3$ , n(k) still increases anomalously for  $k \approx k_{\text{max}}$ , where  $k_{\text{max}}$  is the maximum value of k of given networks with size N. We analytically show that n(k) satisfies the relation  $n(k) \sim k^2 P(k)$  for any degree distribution P(k) and the total number of capture events  $N_{\text{tot}}$  is proportional to  $\langle k^2 \rangle$ , which causes the  $\gamma$ -dependent behavior of S(N, t) and  $\langle T \rangle$ .

DOI: 10.1103/PhysRevE.74.046118

PACS number(s): 89.75.Hc, 05.40.Fb, 89.75.Fb

#### I. INTRODUCTION

The diffusion-controlled reactions, in which diffusing particles are immediately converted to a product if a pair of them meets together, have many physical applications. Examples include electron trapping and recombination, exciton fusion, wetting, melting, and commensurate-incommensurate transitions [1–7]. Among these examples, dynamical properties of wetting, melting, and commensurate-incommensurate transition are known to be related to the diffusive capture process [3–7], whose kinetics can be simplified by diffusing prey-predator model [8,9].

In general, the dynamical properties of diffusioncontrolled reactions are closely related to the first-passage phenomenon of random walks (RWs), because the reaction occurs when diffusing particles first meet a trap or an absorbing boundary [7-9]. Especially, the diffusive capture process can be mapped into coupled RWs [10], in which *n* predator random walkers (lions or moving trap) stalk a prey random walker (a lamb or a moving particle). Since the first-passage probability can be easily obtained from the survival probability [7], the survival probability of a diffusing particle or a lamb is important to understand long-time characteristics of the diffusive capture process [8,9]. The survival probability of a lamb has been well investigated in a *d*-dimensional hypercubic lattice. The diffusive capture process in the spatial dimension d>2 is unsuccessful and the survival probability is finite for any n. In d=2, the capture is successful but not efficient, i.e., the lifetime of a prey random walker is infinite and the survival probability shows a logarithmic decay [11]. The predators efficiently capture a prey for spatial dimension d < 2. On a d-dimensional regular lattice, the survival probability S(t) depends on d as follows [8–10,12,13]:

$$S(t) \sim \begin{cases} t^{-\beta} & \text{for } d < 2, \\ (\ln t)^{-n} & \text{for } d = 2, \\ \text{finite} & \text{for } d > 2, \end{cases}$$
(1)

where the exponent  $\beta$  varies as *n* and the ratio of diffusion constant of walks.

There has been increasing interest in complex networks, since many structures of physically interacting systems are shown to form nontrivial complex networks [14-19]. The early studies on complex networks focused mainly on their topological properties. Recently, the physical systems whose elements interact along the links in complex networks and the effects of the underlying topologies on the dynamical properties of the system have drawn much attention. Many dynamical systems on complex networks show rich behaviors that are far from the mean-field expectations [20-23]. Recently, RWs on complex networks have been studied to show that the dynamical properties of RWs are closely related to the topology of underlying networks [24,25]. This implies that the first-passage property of the diffusive capture process is also affected by the topology of the underlying network. Moreover, several studies on RWs have provided very interesting applications to searching information over the Internet [26,27]. For example, major search engines such as Google and Yahoo use general random walking robots along the links between hypertexts to collect information on each web page [28]. Some peer-topeer (P2P) network protocols also exploit the dynamical properties of RWs on complex networks [26]. Especially, the searching algorithm for a P2P network can be mapped to the system of a diffusing particle to find an immobile absorbing particle (static trap). Our recent study [27] also found that the searching algorithm inspired by the diffusive capture process can dramatically reduce the traffic over the Internet with scalable searching time. The diffusive capture process can also be applied to investigating the condensed phase of the conserved-mass aggregation model on scale-free networks (SFNs) with  $2 < \gamma \leq 3$  [29]. Therefore, it is theoretically and technologically very interesting to study the effect of

<sup>\*</sup>Electronic address: syook@khu.ac.kr

<sup>&</sup>lt;sup>†</sup>Electronic address: ykim@khu.ac.kr

nontrivial complex networks on the diffusive capture process.

In order to understand the relationship between the nontrivial topology of the underlying network and the diffusive capture process on it, we study the random walking lambs and lions on SFNs. The SFNs are the networks whose degree distribution follows  $P(k) \sim k^{-\gamma}$ . The degree exponent  $\gamma$  of the most natural and technological networks is in the range  $2 < \gamma \leq 3$  [15,30]. Moreover, the probability that a random walker visits a node has been shown to be proportional to the degree k of the node [24]. So, the probability that two random walking particles meet at a node is proportional to  $k^2$ . Therefore, we expect that the dynamical properties of the diffusive capture process on the SFNs with the degree exponent  $2 < \gamma \leq 3$  are different from those on the SFNs with  $\gamma > 3$ , because  $\langle k^2 \rangle$  diverges for  $\gamma < 3$ . To verify our expectations, we study the diffusive capture process on the SFNs. More recently, Noh et al. [25] also showed that the return probability of a random walker on the looped SFNs (LSFNs) is completely different from that on the tree-structured scalefree networks (TSFNs). Therefore, the diffusive capture process on TSFNs is also investigated.

This paper is organized as follows. In Sec. II, we present a method of generating TSFNs and dynamical properties of the diffusive capture process on TSFNs. The results on LSFNs are shown in Sec. III. In Sec. IV, we discuss the relation between degrees of SFNs and capture events. Finally, we summarize our results in Sec. V.

#### **II. DIFFUSIVE CAPTURE PROCESS ON TSFNs**

All the models are defined on a graph  $G = \{N, K\}$ , where N is the number of nodes in the networks and K is the number of degrees with the average degree  $\langle k \rangle = 2K/N$ . Networks with the sizes  $N = 10^3 - 10^6$  are generated for each  $\gamma$ .

To make a tree structure, the total number of links in networks with the size N is constrained to be  $K_T=2(N-1)$ . So we use the following method to make TSFNs. Select Ninteger random numbers  $\{k_1, k_2, \dots, k_N\}$  whose probability distribution satisfies  $f(k) \sim k^{-\gamma}$ . Let  $K_0 = \sum_{i=1}^N k_i$ . If  $K_T > K_0$  $(K_T < K_0)$ , then choose  $k_i$  randomly from the set  $\{k_i\}$  and decrease (increase) by 1 as  $k_i \rightarrow k_i + 1$  ( $k_i \rightarrow k_i - 1$ ). If the resulting  $k_i$  becomes 0, the above process is rejected and take the same process again. This process is repeated until  $K_T$  $=K_0$  is fulfilled. Sort the final set  $\{k_1, k_2, \dots, k_N\}$  with  $K_T$  $=\sum_{i=1}^{N} k_i$  to satisfy  $k_1 > k_2 > \cdots > k_N$  and set  $k_i$  to be the degree of the node *i*. Now link the nodes 2, 3,...,  $k_1+1$  to the node 1 to make the degree of the node 1  $k_1$ . Next, link the nodes  $k_1+2, k_1+3, \ldots, k_1+k_2$  to the node 2 to make the degree of the node 2  $k_2$ . Repeat the same processes until no isolated node remains.

Figures 1(a) and 1(b) show snapshots of typical networks generated by the above procedure with 1000 nodes. Figure 1(c) shows the degree distribution of TSFNs constructed with  $\gamma$ =2.4 and 4.3. To generate TSFNs, we use N=10<sup>5</sup>. The data show a good agreement between the exponents  $\gamma$  obtained from the degree distribution,  $P(k) \sim k^{-\gamma}$ , of the constructed network (symbols) and  $\gamma$  of the original distribution



FIG. 1. Schematic structures of TSFNs with  $\gamma$ =4.3 (a) and 2.4 (b) constructed by the method introduced in Sec. II. The plot of the degree distribution of the TSFNs with  $\gamma$ =2.4 and 4.3 (c). Network size is N=10<sup>5</sup>. The degree distribution  $P(k) \sim k^{-\gamma}$  of the constructed networks follows the original distribution  $f(k) \sim k^{-\gamma}$  well.

 $f(k) \sim k^{-\gamma}$  (solid and dashed lines). Figures 1(a) and 1(b) show that the topological properties of TSFNs with  $\gamma > 3$  are completely different from that with  $\gamma < 3$ . For example, the average diameter  $\overline{d}$  scales  $\overline{d} \sim \ln N$  for  $\gamma > 3$  and  $\overline{d} \sim \ln \ln N$ for  $\gamma < 3$  [32]. We find relatively small degree fluctuation in the networks with  $\gamma > 3$  [see Fig. 1(a)]. The small fluctuation causes almost homogeneous topological properties. Thus, we expect that the dynamical properties on the networks with  $\gamma > 3$  follow the mean-field expectations, i.e., the nonzero survival probabilities of a lamb in the  $N \rightarrow \infty$  limit. However, for a network with  $\gamma < 3$ , there exist many short branches and a big hub having most links. The dynamical properties of a given system on this kind of network can be dominated by this hub. Therefore, as we shall show later, the lamb can easily hide from the lions by moving to a different branch through the hub.

Now, let us discuss the dynamical properties of the diffusive capture process on the constructed TSFNs. Initially the lamb and *n* lions are randomly distributed to the nodes on the networks. In the simulations, we consider two updating rules to determine the position of the random walking lamb and lions: (i) parallel update and (ii) random update. In the parallel update rule, we move all lions and a lamb at the same time, then check if the lamb meets the lion(s). In this algorithm, the lamb and the lions should be initially separated by 2m (m=1,2,...) steps. Otherwise, the lamb never meets the lion moving on TSFNs. In the random update rule, we choose one among the lamb or lions at random. The randomly chosen lamb or lion takes a random walk and then we check if the lamb meets the lion. The unit time is defined by (n+1) such random movements. We find that there is no significant difference between dynamical behaviors for each update rule. We also find that the dynamical behaviors are not varied critically if n is varied [see Figs. 2(a) and 2(b)]. Therefore, we will mainly show the simulation results for



FIG. 2. Average lifetime  $\langle T \rangle$  on TSFNs for (a) n=1 and (b) n=10 with various  $\gamma$  (=4.3, 3.0, and 2.4). The lines represent the best fits of  $\langle T \rangle \sim N^{\alpha}$  to data. The obtained exponents are  $\alpha = 1.1 \pm 0.1$  (solid line) for  $\gamma = 4.3$ ,  $\alpha = 1.0 \pm 0.1$  (dashed line) for  $\gamma = 3.0$ , and  $\alpha = 0.7 \pm 0.1$  (dotted line) for  $\gamma = 2.4$ . For  $\gamma = 2.4 \langle T \rangle$  substantially deviates from the  $\langle T \rangle \sim N^{0.7}$  line when  $N = 10^6$ .

one lamb and ten lions in this paper. All the numerical results are obtained by averaging over 100 network realizations and 1000 histories for each network realization.

Figure 2 shows the plot of the average lifetime  $\langle T \rangle$  on TSFNs with various  $\gamma$  against N using the random update rule.  $\langle T \rangle$  increases with N for any  $\gamma$  and n. The lines represent the best fits of the relation  $\langle T \rangle \sim N^{\alpha}$  to the data. From the fits, we find that  $\langle T \rangle$  follows the power law  $\langle T \rangle \sim N^{\alpha}$  with  $\alpha = 1.1 \pm 0.1$ ,  $1.0 \pm 0.1$ , and  $0.7 \pm 0.1$  for  $\gamma = 4.3$  (solid lines), 3.0 (dashed lines), and 2.4 (dotted lines), respectively. But, as shown in Fig. 2,  $\langle T \rangle$  for  $\gamma = 2.4$  deviates substantially from the power-law relation  $\langle T \rangle \sim N^{0.7}$  as increasing N (square). By considering the error bars of the  $\alpha$ , we can conclude that  $\alpha \simeq 1.0$  for  $\gamma \ge 3$  and  $\alpha \simeq 0.7$  for  $\gamma < 3$ . The data in Figs. 2(a) and 2(b) indicate that the values of  $\alpha$  do not depend on n. The inset of Fig. 2(b) shows the results obtained by using the parallel update rule. The data obtained from the parallel update rule show nearly the same dynamical behaviors as those obtained from the random update rule.

We also measure the survival probability on TSFNs with various  $\gamma$  and N. Figure 3 shows S(N,t) for  $\gamma=4.3$  (a), 3.5 (b), 3.0 (c), 2.75 (d), and 2.4 (e) from the random update rule. As is shown in Fig. 3, the survival probability for

![](_page_2_Figure_7.jpeg)

FIG. 3. The scaled plots of the survival probability S(N,t) on TSFNs with  $\gamma$ =4.3 (a), 3.5 (b), 3.0 (c) and the plot of S(N,t) for  $\gamma$ =2.75 (d) and  $\gamma$ =2.4 (e). For the scaling plots,  $\alpha \approx 1$ , which is obtained from the fits for  $\gamma > 3$  in Fig. 2, is used. The insets show the same S(N,t) in log-log scale.

 $\gamma \ge 3$  decays exponentially as  $S(N,t) \sim \exp[-t/\tau(N)]$ , however S(N,t) for  $\gamma < 3$  manifests a substantial deviation from an exponential decay. If we assume that  $S(N,t)=S_oe^{-t/\tau}$ , then we obtain  $\langle T \rangle \sim \tau(N)$  from the relation  $\langle T \rangle$ 

![](_page_3_Figure_1.jpeg)

FIG. 4. Plot of N(L,T) on TSFN with  $\gamma=2.4$  and  $N=10^4$ . The initial separation *l* of the lamb and the lion is 2. *L* is the maximum separation between the lamb and the lion during random walks (symbols).

=  $\int_{0}^{\infty} t(-dS(N,t)/dt)dt$ . Therefore, the characteristic time  $\tau$  of S(N,t) scales in the same manner as  $\langle T \rangle$ , and the scaling t/N collapses S(N,t)'s for various N into one curve when  $\gamma \ge 3$  [see Figs. 3(a)-3(c)]. The relation  $\tau(N) \sim N$  implies that S(N,t) becomes a nonzero constant in the limit  $N \rightarrow \infty$ . Therefore, the exponentially decaying S(N,t) for  $\gamma > 3$  can be regarded as the finite-size effect. However, S(N,t) for TSFN with  $\gamma < 3$  shows the long-living tail [Figs. 3(d) and 3(e)], which decreases more slowly than exponential but faster than power-law decay (see the insets). Hence, we expect

$$\lim_{t \to \infty} \lim_{N \to \infty} S(N, t) = \begin{cases} \text{long-living tail} & (\gamma < 3), \\ \text{finite} & (\gamma \ge 3). \end{cases}$$
(2)

The origin of  $\gamma$ -dependent behaviors of S(N,t) will be discussed in Sec. IV.

We now discuss the origin of the long-living tail of S(N, t)for  $\gamma = 2.4$ . One possible explanation of the slowly decaying tail of S(N,t) can be understood from the topological properties of TSFNs. As is shown in Fig. 1(b), the biggest hub in the TSFN with  $\gamma < 3$  has many short branches. Hence, if the lamb hides from the lions by moving into a different branch, then the lion should visit almost all branches to search for a hidden lamb. Thus, the hidden lamb can live longer than the one in the same branch with the lion. Therefore, S(N,t) on TSFN with  $\gamma < 3$  can have the long-living tail. In order to verify this argument, we count the number of capture events N(L,T) with the lifetime T among 10<sup>6</sup> different histories on TSFN with  $\gamma = 2.4$  (Fig. 4) by using the random update rule. At the same time, we also measure the maximum separation L between the lamb and the lion during random walks. Initially, we set the shortest path length *l* between the lamb and lion to be 2. We use the TSFN whose average diameter  $\overline{d}$  is  $\overline{d} \simeq 4$ . If  $L \simeq \overline{d}$ , then the lamb and the lion have once been in different branches during random walks. The data in Fig. 4 explicitly show that N(L,t) for L=2 decays fast (the maxi-

![](_page_3_Figure_8.jpeg)

FIG. 5. The scaled plots of the survival probability S(N,t) on LSFNs with  $\gamma$ =4.3 (a), 3.0 (b), and 2.4 (c). To scale the time *t*,  $\alpha$ 's are obtained from the fits of  $\langle T \rangle \sim N^{\alpha}$  to the data (which are not shown in this paper).

mum lifetime of a lamb is  $T_{\text{max}} \approx 10^2$ ). In contrast, N(L,t) for L=4 shows the long-living tail up to  $T_{\text{max}} \approx 10^4$ . From the relation  $S(t) \propto -\int_t \Sigma_L N(L,t) dt$ , one can easily see that the long-living tail of S(t) is caused by the diffusive capture process with L=4.

# **III. DIFFUSIVE CAPTURE PROCESS ON LSFNs**

The dynamical behavior of the one lamb and one lion model with the random update rule on LSFNs has already been reported in Ref. [29]. We measure the S(N,t) and  $\langle T \rangle$ for various *n*'s as the random and the parallel update algorithms. We check that  $\langle T \rangle$  and S(N,t) scale in the same manner as those reported in Ref. [29]. Therefore, we will report our results briefly in this section. To generate the LSFNs of sizes  $N=10^3-10^6$ , we use the static network model suggested by Ref. [31] with  $\langle k \rangle = 4$ . In Fig. 5, we display S(N,t)for n=10 on LSFNs with various  $\gamma$  (=4.3, 3.0, and 2.4). S(N,t) exponentially decays as  $S(N,t) \sim \exp(-t/\tau)$  for all  $\gamma$ 's. As discussed in Sec. II, by assuming that S(N,t) $\sim \exp(-t/\tau)$ , we obtain the relation  $\langle T \rangle \sim \tau(N)$ . Using the value of  $\alpha$ 's obtained from the best fit of  $\langle T \rangle \sim N^{\alpha}$  to the data (which are not shown), we scale the time as t/N for  $\gamma > 3$ 

![](_page_4_Figure_1.jpeg)

FIG. 6. Plot of capture event distribution n(k) on LSFN (a) and TSFN (b). The lines stand for the relation  $n(k) \sim k^{-\sigma}$ . The solid lines correspond to  $\sigma$ =0.4 and the dotted lines represent  $\sigma$ =2.3. The data for  $k < 10^3$  (a) and for  $10^2 < k < 10^3$  (b) show good agreement with the relation  $\sigma = \gamma - 2$ . The dashed lines correspond to  $\sigma$ =-1.35 (a) and  $\sigma$ =-1.9 (b) representing the anomalous increases of n(k)'s for large k.

and as  $t/N^{0.55}$  for  $\gamma=2.4$ . As shown in Fig. 5(a), S(N,t) for  $\gamma>3$  collapses well into a single curve. This result implies that  $\tau(N) \sim N$  for  $\gamma>3$ . However, as we can see in Fig. 5(c), S(N,t) for  $\gamma=2.4$  does not collapse into a single curve by the simple scaling. Note that  $\gamma=3$  is marginal, below which  $\langle k^2 \rangle$  diverges. Thus, the logarithmic correction for the scaling of  $\tau$  can be expected for  $\gamma=3$ . In Fig. 5(b), we scale the time as  $t/N(\log_{10}N)^{-\delta}$  with  $\delta=1.6\pm0.2$  and find that S(N,t) for  $\gamma=3$  collapses into one curve. Here  $\delta$  is obtained by the best fit of  $\langle T \rangle$  to the function  $\langle T \rangle \sim N(\log_{10}N)^{-\delta}$ . This means that  $\tau$  scales as  $\tau \sim N(\log_{10}N)^{-1.6}$ . Therefore, we expect

$$\lim_{t \to \infty} \lim_{N \to \infty} S(N, t) = \begin{cases} S_o e^{-t/\tau_\infty} & (\gamma < 3), \\ \text{finite} & (\gamma \ge 3), \end{cases}$$
(3)

where  $\tau_{\infty}$  is the asymptotic value of  $\tau$ , and depends on  $\gamma$ . Equations (2) and (3) show one of the pieces of evidence that the dynamical behavior of the diffusive capture process is related to  $\langle k^2 \rangle$ .

### IV. RELATION BETWEEN DEGREES AND CAPTURE EVENTS

In the previous sections, we showed that the underlying topology significantly affects the dynamical behaviors of S(N,t) and  $\langle T \rangle$  of a lamb. For a quantitative description of the relationship between the structure of the networks and the dynamical properties of the lamb-lion model, we count the capture events occurring at the nodes with degree k. Figure 6(a) shows the normalized number of capture events n(k) on LSFNs with one lamb and ten lions. As we can see, n(k)'s both on LSFNs and TSFNs with  $\gamma$ =4.3 decay as  $n(k) \sim k^{-\sigma}$  and the obtained  $\sigma$  is  $\sigma \approx 2.3$ . However, n(k) measured on LSFNs with  $\gamma$ =2.4 decreases with  $\sigma \approx 0.4$  for  $k < 10^3$  and

increases with  $\sigma \approx -1.35$  for  $k > 10^3$ . On TSFNs with  $\gamma = 2.4$ , we also find  $\sigma \approx 0.4$  for  $10^2 < k < 10^3$  and  $\sigma \approx -1.9$  for  $k > 10^3$  [see Fig. 6(b)]. The decreasing behavior of n(k) for moderate k can be explained by the following analysis. In the steady state, the probability that a random walker visits a node *i* with degree  $k_i$  is given by [24]

$$P_i^{\infty} = Ak_i, \tag{4}$$

where  $A=1/\sum_i k_i$ . Therefore, the normalized capture events simply can be written as

$$n(k) = \sum_{i,j=1}^{N} P_i^{\infty} \delta_{k_i,k} P_j^{\infty} \delta_{i,j} = A^2 k^2 N P(k).$$
 (5)

If  $P(k) \sim k^{-\gamma}$ , then we obtain

$$\sigma = \gamma - 2. \tag{6}$$

Since we do not assume any topological property of underlying networks, Eq. (6) is valid for any SFNs and the obtained values of  $\sigma$  from the simulations for  $k < 10^3$  are consistent with Eq. (6). Note that the n(k)'s anomalously increase for  $k > 10^3$  both on the LSFNs and the TSFNs with  $\gamma=2.4$ . This anomalous behavior of n(k) can be explained by the rarity of the dominant hubs in complex networks, i.e., the P(k) for  $k \approx k_{\text{max}}$  does not change much as k varies. Here  $k_{\text{max}}$  is the degree of the biggest hub. Thus, if we assume that  $P(k \approx k_{\text{max}}) \sim \text{const}$ , then n(k) increases as  $n(k) \sim k^2$  or

$$\sigma = -2 \tag{7}$$

for  $k \approx k_{\text{max}}$ . The obtained  $\sigma(k \approx k_{\text{max}})$ 's for  $\gamma = 2.4$  from the numerical simulations are -1.35 on LSFN and -1.9 on TSFN. The value of  $\sigma(k \approx k_{\text{max}})$  obtained from the simulation on LSFN with  $\gamma = 2.4$  is considerably smaller than Eq. (7). This disagreement might originate from our assumption that  $P(k \approx k_{\text{max}})$  does not depend on k. For example, if we assume that  $P(k \approx k_{\text{max}}) \sim k^{-\epsilon}$  with  $\epsilon < 2$ , then  $\sigma$  becomes  $\sigma = \epsilon - 2(<0)$ . Equations (5) and (7) thus provide the topological origin of the gathering behavior of random walkers at the dominant hubs. This implies that the lamb and the lion have a strong tendency to move into big hubs, i.e., the hubs collect the lamb and the lion. Therefore, the hubs in the networks with strong heterogeneity become the effective attractors of random walking lambs and lions.

Another important consequence of Eq. (5) is that the dynamical properties of the lamb-lion model depend on  $\langle k^2 \rangle$ . Note that from Eq. (5), the total number of capture event  $N_{\text{tot}}$ can be expressed as

$$N_{\text{tot}} \propto \sum_{k} n(k) \sim \sum_{k} k^2 P(k) \sim \langle k^2 \rangle.$$
(8)

If  $\gamma < 3$ , then  $N_{\text{tot}}$  also diverges. This means that the lions on the networks can capture the lamb no matter how the lamb is trying to escape from the lions. Therefore,  $S(N,t) \rightarrow 0$  in the limit  $N \rightarrow \infty$ . For  $\gamma > 3$ ,  $N_{\text{tot}}$  converges to a finite value and the lions can catch the lamb with a probability p < 1, thus  $S(N,t) \neq 0$  for large N. The expectations agree well with our simulation results shown in Secs. II and III. Based on Eq. (8), we also expect that the long-living tail of S(N,t) on TSFN with  $\gamma < 3$  eventually goes to 0. Moreover, if we increase the number of lions, *n*, then  $N_{\text{tot}}$  or n(k) increases by *n* times. Therefore, the scaling behavior of  $\langle T \rangle$  is not affected by *n*, which is confirmed by numerical simulation (see Fig. 2). Thus, we can expect that the scaling behavior of S(N,t) does not depend on *n*.

### V. SUMMARY AND DISCUSSION

We investigate the diffusive capture process on two different kinds of SFNs. The diffusive capture process can be mapped into the diffusing lamb and lion model. We measure S(N,t) and  $\langle T \rangle$  of a lamb. We find that on both TSFNs and LSFNs, S(N,t) of a lamb decays exponentially for  $\gamma > 3$ . Simple scaling arguments predict that the exponentially decaying S(N,t) for  $\gamma > 3$  originates from the finite-size effect. This result recovers the mean-field expectation S(N,t) $\rightarrow$  finite in the limit  $N \rightarrow \infty$ . However, we find that S(N,t) on TSFNs with  $\gamma < 3$  has the long-living tail. By measuring the number of capture events with lifetime *T*, we find that the long-living tail is caused by enormously many short branches in TSFNs. On LSFN with  $\gamma < 3$ , we find that S(N,t)decays exponentially and no lamb can survive in the thermodynamic limit.

The finite survival probability for  $\gamma > 3$  can be easily understood by analogy with a random tree or high-dimensional regular lattice. If  $\gamma$  is sufficiently large, then the network becomes homogeneous and its structure can be well approximated by a random tree or high-dimensional lattice on which S(N,t) is known to be finite [see Eq. (1)]. We also find that

 $\gamma$ =3 is marginal and below which S(N,t) significantly deviates from the high-dimensional expectations.

In order to understand the non-mean-field behavior of S(N,t) for  $\gamma < 3$ , we study the normalized number of capture events n(k) both numerically and analytically. From the numerical simulations, we find that n(k) decreases as a power law,  $n(k) \sim k^{-\sigma}$ . The simple analytic derivation of n(k) provides the relation  $\sigma = \gamma - 2$ , which shows a good agreement with numerical results for moderate k. However, we find that n(k) anomalously increases for large k in the simulations. This means that most of the lambs are captured by lions at the hubs. Hence, those hubs become the effective attractors. By assuming the rarity of dominant hubs and  $P(k) \approx \text{const}$ for large  $k(\approx k_{\text{max}})$ , we show that the n(k) can increase for large k. From the analytic expectation of n(k), we also show that  $N_{\rm tot} \sim \langle k^2 \rangle$ . This implies that if the fluctuation in degree is large enough ( $\gamma < 3$ ), then the dynamical behaviors of the diffusive capture process deviate from the mean-field expectations by emerging effective attractors of random walkers. From the analytic expectation of n(k) and  $N_{tot}$ , we also expect that the long-living tail of S(N,t) observed on TSFN with  $\gamma < 3$  [see Eq. (2)] will eventually go to 0 due to the diverging  $\langle k^2 \rangle$ .

#### ACKNOWLEDGMENTS

We thank to Dr. Kwon and Dr. Yoon for useful discussions. This work is supported by Korea Research Foundation Grant No. KRF-2004-015-C00185.

- M. Agranovich and M. D. Galanin, *Electronic Energy Transfer* in Condensed Matter (North-Holland, Amsterdam, 1985).
- [2] S. J. Parus and R. Kopelman, Phys. Rev. B 39, 889 (1989).
- [3] D. ben-Avraham, J. Chem. Phys. 88, 941 (1988).
- [4] M. E. Fisher, J. Stat. Phys. 34, 667 (1984).
- [5] D. A. Huse and M. E. Fisher Phys. Rev. B 29, 239 (1984).
- [6] R. Lipowsky and T. M. Nieuwenhuizen, J. Phys. A 21, L89 (1988).
- [7] S. Redner, A Guide to First-Passage Processes (Cambridge University Press, New York, 2001).
- [8] P. L. Krapivsky and S. Redner, J. Phys. A 29, 5347 (1996).
- [9] S. Redner and P. L. Krapivsky, Am. J. Phys. 67, 1277 (1999).
- [10] M. Bramson and D. Griffeath, *Capture Problems for Coupled Random Walks*, in Random Walks, Brownian Motion, and Interacting Particle Systems: A Festschrift in Honor of Frank Spitzer, edited by R. Durrett and H. Kesten (Birkhauser, Boston, 1991), pp. 153–188.
- [11] G. H. Weiss, *Aspects and Applications of the Random Walk* (North-Holland, Amsterdam, 1994).
- [12] W. Feller, An Introduction to Probability Theory (Wiley, New York, 1971), Vol. 1.
- [13] Aspects and Applications of the Random Walk (North-Holland, Amsterdam, 1994).
- [14] H. Jeong, B. Tombor, R. Albert, Z. N. Oltvai, and A.-L. Barabási, Nature (London) 407, 651 (2001); S.-H. Yook, Z. N.

Oltvai, and A.-L. Barabási, Proteomics 4, 928 (2004).

- [15] S. N. Dorogovtsev and J. F. F. Mendes, Adv. Phys. **51**, 1079 (2002); R. Albert and A.-L. Barabási, Rev. Mod. Phys. **74**, 47 (2002).
- [16] R. Albert, H. Jeong, and A.-L. Barabási, Nature (London) 401, 130 (1999).
- [17] A. Broder, R. Kumar, F. Maghoul, P. Raghavan, S. Rajalopagan, R. Stata, A. Tomkins, and J. Wiener, Comput. Netw. 33, 309 (2000).
- [18] M. Faloutsos, P. Faloutsos, and C. Faloutsos, Comput. Commun. Rev. 29, 251 (1999).
- [19] S.-H. Yook, H. Jeong, and A.-L. Barabási, Proc. Natl. Acad. Sci. U.S.A. 99, 13382 (2002).
- [20] A. V. Goltsev, S. N. Dorogovtsev, and J. F. F. Mendes, Phys. Rev. E 67, 026123 (2003).
- [21] M. Leone, A. Vázques, and R. Zecchina, Eur. Phys. J. B 28, 191 (2002).
- [22] R. Cohen, K. Erez, D. ben-Avraham, and S. Havlin, Phys. Rev. Lett. 85, 4626 (2000).
- [23] R. Cohen, D. ben-Avraham, and S. Havlin, Phys. Rev. E 66, 036113 (2002).
- [24] J. D. Noh and H. Rieger, Phys. Rev. Lett. 92, 118701 (2004).
- [25] J. D. Noh and S.-W. Kim, J. Korean Phys. Soc. 48, S202 (2006).
- [26] L. A. Adamic, R. M. Lukose, A. R. Puniyani, and B. A. Hu-

berman, Phys. Rev. E 64, 046135 (2001).

- [27] S. Lee, S.-H. Yook, and Y. Kim (unpublished).
- [28] R. Pestor-Satorras and A. Vespignani, *Evolution and Structure* of the Internet: A Statistical Physics Approach (Cambridge University Press, Cambridge, 2004).
- [29] S. Kwon, S. Lee, and Y. Kim, Phys. Rev. E 73, 056102 (2006).
- [30] L. A. Adamic and B. A. Huberman, Science **287**, 2115 (2000).
- [31] K.-I. Goh, B. Kahng, and D. Kim, Phys. Rev. Lett. 87, 278701 (2001).
- [32] R. Cohen and S. Havlin, Phys. Rev. Lett. 90, 058701 (2003).